A. White

COULD THE LHC POMERON BE THE KEY TO EVERYTHING?*

A color sextet quark sector of QCD could be the origin of EW symmetry breaking. It would dominate UHE x-sections & would explain the CR knee, dark matter, & much more.

It should be unavoidably seen by FP420 detectors via pair production of W's &Z's. In this talk I will raise the stakes by describing why

THE SEXTET SECTOR MUST BE PART OF QUD - A UNIQUE SU(5)
UNIFIED MASSLESS "THEORY OF MATTER" (no gravity) THAT COULD
BE THE ORIGIN OF THE FULL STANDARD MODEL.

Discovery of the sextet sector implies the discovery of QUD !!!

^{*} Presented at - Physics with FP420 - Manchester, December 2006.

Key points of my previous talks at this meeting were

- 1. The Critical Pomeron $\{$ the only known solution of high-energy unitarity $\}$ appears in QCD when the asymptotic freedom constraint is saturated by 6 color triplet quarks + 2 color sextet quarks \rightarrow "QCD $_S$ "
- 2. $W^{\pm} \& Z^{0}$ will eat the "sextet pions" \longrightarrow electroweak symmetry breaking with no new interaction (the electroweak scale is the QCD sextet chiral scale !!!)
- 3. The CR knee, & other CR phenomena, suggest a threshold for the expected new physics* between Tevatron & LHC. If so, there must be large x-section effects at the LHC.
- 4. The double \mathbb{P} x-section will be crucial for establishing that the sextet sector has appeared.

BUT, how is the electroweak anomaly of the sextet sector canceled (??) & how are other particle masses generated ???

^{*} ARW - Phys. Rev. D72:036007 (2005).

Some years ago, Kang & I looked at left-handed unified theories and found a remarkable, *but puzzling*, result. *Requiring* only

- 1. the sextet sector
- 2. asymptotic freedom
- 3. anomaly cancelation

uniquely selects SU(5) gauge theory with the fermion representation $5 + 15 + 40 + 45^* \longleftrightarrow$

"QUD"

 $\{Quantum\ Uno/Unification/Unitary/Underlying\ Dynamics\}$

Amazingly, the triplet quark and lepton sectors of QUD (which were not asked for !) are very close to the Standard Model !!

There are 3 "generations" of quarks & antiquarks with charges $\pm 2/3, \ \pm 1/3$ (and so QUD contains QCD $_S$) together with 3 "generations" of leptons.

Under $SU(3)\otimes SU(2)\otimes U(1)$

$$\begin{split} \mathbf{5} &= (\mathbf{3}, \mathbf{1}, -\frac{1}{3}))^{\{\mathbf{3}\}} + (\mathbf{1}, \mathbf{2}, \frac{1}{2}))^{\{\mathbf{2}\}} \;, \\ \mathbf{15} &= (\mathbf{1}, \mathbf{3}, \mathbf{1}) + (\mathbf{3}, \mathbf{2}, \frac{1}{6}))^{\{\mathbf{1}\}} + (\mathbf{6}, \mathbf{1}, -\frac{2}{3}) \;, \\ \mathbf{40} &= (\mathbf{1}, \mathbf{2}, -\frac{3}{2}))^{\{\mathbf{3}\}} + (\mathbf{3}, \mathbf{2}, \frac{1}{6})^{\{\mathbf{2}\}} + (\mathbf{3}^*, \mathbf{1}, -\frac{2}{3}) + (\mathbf{3}^*, \mathbf{3}, -\frac{2}{3}) \\ &\quad + (\mathbf{6}^*, \mathbf{2}, \frac{1}{6}) + (\mathbf{8}, \mathbf{1}, \mathbf{1}) \;, \\ \mathbf{45}^* &= (\mathbf{1}, \mathbf{2}, -\frac{1}{2}))^{\{\mathbf{1}\}} + (\mathbf{3}^*, \mathbf{1}, \frac{1}{3}) + (\mathbf{3}^*, \mathbf{3}, \frac{1}{3}) + (\mathbf{3}, \mathbf{1}, -\frac{4}{3}) \\ &\quad + (\mathbf{3}, \mathbf{2}, \frac{7}{6}))^{\{\mathbf{3}\}} + (\mathbf{6}, \mathbf{1}, \frac{1}{3}) + (\mathbf{8}, \mathbf{2}, -\frac{1}{2}) \end{split}$$

The "generations" $\{1\},\{2\},\{3\}$, are scattered amongst the separate representations. It will be very important later that the complete representation is a vector theory wrt SU(3)xU(1) $_{em}$.

But, the SU(2)xU(1) quantum numbers are not quite right. Also, there are (unwanted?) color octet quarks with lepton-like electroweak quantum numbers.

It was very frustrating that we had found a unique theory that is almost, but not quite, the Standard Model!!

- As massless field theories QUD & QCD_S have similar, very special, UV
 & IR properties.
 - => the high-energy bound-state S-Matrix can be constructed via multiregge theory & infra-red chiral anomalies.
- Only after I understood the physics of massless QCD_S , did it become apparent to me what the true role of QUD could be.

I realized, incredibly, that QUD could be to the full Standard Model what QCD is to the hadronic sector !!!! In the QUD bound-state S-Matrix

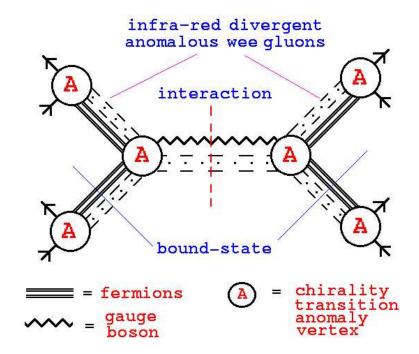
SU(5) COLOR is CONFINED, not just SU(3) color, & so ALL ELEMENTARY GAUGE BOSONS AND FERMIONS ARE CONFINED - AND MASSLESS.

- All hadrons and leptons would be QUD bound-states and all interactions would be composite. Is this possible ???
- Because QUD is real wrt $SU(3)xU(1)_{em}$ and the lepton and triplet quark sectors are so close to the Standard Model, it could be !!

Very briefly -

we construct bound states and interactions similarly in QCD_S and QUD, i.e. via multi-regge amplitudes that contain **infra-red divergent** gauge bosons **coupled to anomalies**.

Restoring the gauge symmetry in steps & extracting anomaly infra-red divergences -



bound-states appear as Goldtsone boson "anomaly poles" formed as color zero combinations of fermions in an "anomalous wee gluon" background *



interactions are color zero combinations of a finite transverse momentum gauge boson in the same wee gluon background.



^{*}Wee gluons \equiv some fermions are in negative energy states

- To obtain the states & amplitudes of QCD_S, we start in "ColorSuperconducting" QCD_S $\{SU(3) \text{ color} \rightarrow SU(2)\}$.
- ullet The physical states of QCD_S are Goldstones in " $CSQCD_S$ ".
- → triplet meson & nucleon states
- → no hybrid sextet/triplet quark states
- \rightarrow sextet "pions" & "nucleons" ($P_6\&N_6$)

Consistent with, but much less than just requiring confinement & chiral symmetry breaking.

The sextet neutron, the N_6 , will be stable & dominate UHE x-sections

→ Dark Matter !!

- In CSQCD_S, the interaction is a massive gluon reggeon in an anomalous wee gluon condensate \leftrightarrow the "supercritical" \mathbb{P} .
 - => In QCD $_S$, the interaction is the Critical ${\Bbb P}$ \leftrightarrow regge pole + interactions
 - => no BFKL pomeron, no odderon, & no glueballs

The QCD_S states are much fewer & the interaction much simpler, than in conventional QCD, in better agreement with experiment !!

For QUD, it is crucial that "anomalous" divergences are exponentiated by left-handed gauge bosons.

=> divergences
→ maximal non-abelian vector subgroup
=> a strong interaction involving only
SU(3) singlet combinations !!

To obtain amplitudes we, again, build up the symmetry in stages

1	SU x	$\frac{(3) \text{ colo}}{x'x'}$	r x"x" \	
	x' x'	SU(2)C	0 0) (1)
	x'' x''	0 0	SU(2) _L	SU(4)

With $SU(2)_C$ (vector) symmetry the states are Goldstone π_C 's, i.e qq, $\bar{q}\bar{q}$, & $q\bar{q}$ pairs in a condensate.

The q's are SU(3) 3's, 6's, & 8's. 8's are real wrt SU(3), but contain complex doublets wrt $SU(2)_C$.

The exchanges producing interactions are

- 1. A massive x gluon in the condensate $\leftrightarrow \mathbb{P}$.
- **2.** $SU(2)_L \otimes U(1)$ bosons in the condensate $\leftrightarrow W^{\pm,0}$, Y.
- 3. A massive x" boson in the condensate will be confined by SU(3) color

Wee gluon interactions give W^\pm & Z^0 a mass \leftrightarrow mixing with π_C 's.

When SU(4) symmetry is restored, "leptons" appear as bound states of elementary leptons and "octet pions", e.g. the e/ν will be $(1,2,-\frac{1}{2})\times(8,1,1)\times(8,2,-\frac{1}{2})\leftrightarrow$ SU(5) singlet $45^*\times40\times45^*$. (The μ & the τ contain 3 elementary leptons.)

When SU(5) symmetry is restored, we anticipate that

- 1. The ℙ becomes critical.
- 2. The wee gluon component of the γ & the W^{\pm}, Z^0 becomes even signature (essentially, the zero k_{\perp} component of the \mathbb{P}).
- 3. The octet π 's are no longer Goldstones & disappear from the low k_{\perp} region.
- 4. Also SU(3) reality => octet π 's have no (anomaly) coupling to the $\mathbb{P}=>$ leptons have no strong interaction and no infra-red SU(3) mass generation.
- 5. Because octet π 's contribute only at large k_{\perp} the $SU(2)_L \otimes U(1)$ symmetry will appear in low k_{\perp} interactions (as $sextet\ SU(2)\ flavor$).
- 6. The $SU(2)_L \otimes U(1)$ anomaly => three generations of "hadrons" and "leptons".
- 7. $SU(2)\otimes U(1)$ quantum numbers of the octet π 's => low k_{\perp} states will have the the singlet/doublet structure of the Standard model.

The octet quarks (which at first sight seemed unwanted) are fundamental for the underlying SU(5) invariance and the generation structure of states.

- Much of what I have described needs to be better established and very many questions remain to be answered.
- Also, what I am proposing is very radical wrt the current theoretical paradigm !!
- I am saying that the Standard Model has an underlying unifying field theory, but, it is massless & has a very small coupling $(\alpha_u < O(1/50) \leftrightarrow an\ IR\ fixed-pt.)$ It is, almost, conformally invariant.
- Mass-scales are generated by reggeization, mixing, & anomaly interactions, but only in the bound-state S-Matrix - which has a very special role.
- There is no Higgs field although the η_6 is analogous to the Higgs wrt EW symmetry breaking.

There are many general features that are encouraging, including

- 1. The experimentally attractive SU(5) value of the Weinberg angle should hold, even though there is no proton decay!
- 2. Small α_u could be the explanation of small neutrino masses.
- 3. The existence and dominance of Dark Matter is naturally explained.
- 4. The high mass QCD sector produces unification without supersymmetry.
- 5. There are no unwanted symmetries constraining the mass spectrum.
- 6. The SO(10) 144 contains QUD relevant for { string?} unification with gravity?

- To suggest that new strong interaction physics beyond the Standard Model is the "key to everything" & that it should be looked for via the \mathbb{P} is very unconventional.
- Of course, the discovery of new large x-section effects at the LHC would soon make the discussion of such physics conventional !!

Persistently, & singularly, searching for the Critical \mathbb{P} , & pushing my specialist knowledge of multi-regge theory to the limits of my understanding (& beyond!), I arrived first at the sextet sector & now, uniquely, at QUD.

Obviously, it would be incredible if the Standard Model, with all of it's complexity, has the underlying simplicity I have suggested. Nevertheless,

all the necessary ingredients are present &, if the predicted effects of the sextet sector are seen at the LHC, interest in QUD will surely rise rapidly!

I will finish by briefly reviewing how double ${\mathbb P}$ processes can provide the proof that a sextet sector has appeared.

What Should be Seen at the LHC?

At high luminosity, major evidence for the sextet sector would be

- Multiple vector boson and jet x-sections that are much, much, larger than expected, producing a dramatic rise in the average $|p_{\perp}|$
 - from the low energy hadron scale towards the electroweak scale.
- But, there will be other explanations black holes, sphalerons, ...

A priori,

- the neutral N_6 { dark matter}, with a best guess mass $\sim 500~GeV$, will be difficult to detect, since missing energies of several hundred GeV will be common.
- ullet The P_6 , assuming it is not too unstable, should be seen.
- Again, a massive, charged, particle with a large production x-section will not be immediately identified with the sextet sector!

- The double P x-section could be the most definitive early evidence for the existence of the sextet sector.
- With the P's detected via Roman pots, the environment is cleaner & more controlled.

W&Z pairs will be produced in the double \mathbb{P} x-section via sextet pion anomaly poles. {As pion pairs dominate the double \mathbb{P} x-section at low mass, so W&Z pair production will dominate the x-section at the EW mass scale.}

- => when $|k_{\perp}|$ is EW scale, the amplitude is comparable with a jet amplitude that has, apart from anomaly loops that are O(1), the same propagators & couplings
- => at large k_{\perp} , double $\mathbb{P}\ W\&Z$ pairs will give jet x-sections that are comparable with the non-diffractive x-sections predicted by standard QCD.

Generally, a factor of $\left[\frac{F_{\pi_6}}{F_{\pi_3}}\right]^4 \left(\gtrsim O(10^{12})\right)$ is involved in relating sextet and triplet sector x-sections.

The central $\{\mathbb{P}W^+W^-\mathbb{P}\}\$ & $\{\mathbb{P}Z^0Z^0\mathbb{P}\}\$ vertices will vary only slowly with k_{\perp} , but the hadron/ \mathbb{P} vertices have strong k_{\perp} -dependence that should give an extremely large x-section at small t.

- In the low luminosity running, this x-section could be detected by TOTEM in combination with the CMS central detector (if it is operational) where it should be straightforward to look for W&Z pairs.
- Some spectacular events would be expected, in which protons are tagged and only (a multitude of) large E_T leptons are seen in the central detector?

A very large double ${\mathbb P}$ x-section for W&Z pairs

- => longitudinal components of W&Z have direct strong interactions
 - => existence of the sextet sector !!!
- **FP420** will take over later & should surely see the enhanced x-section (whether or not it has been seen by CMS/TOTEM) if it is present !!
- Indeed, with the planned parameters for FP420, the sextet W&Z pair x-section will overwhelm all other physics.

After the combination of \mathbb{P} , W/Z, & jet physics has established that sextet quark physics is definitively discovered, the search for "Dark Matter" will become all important.

The x-section for double ${\mathbb P}$ production of $\{{\sf stable}\}$ $N_6\bar{N}_6$ pairs (with mass $\gtrsim 1~TeV$) could be large enough that it will be definitively seen by an optimum combination of forward pots. It will be a spectacular process to look for.

- The tagged protons determine a very massive state was produced.
- No charged particles seen in any of the detectors.
- ullet Having low energy, the N_6 hadronic x-section will, probably, be small but some hadronic activity may be seen in the central calorimeter
- Charged lepton comparison would allow a separation wrt the multiple Z production of neutrinos.

If the P_6 is relatively stable, & not to different in mass, it would be much simpler to first detect $P_6\bar{P}_6$ pairs